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Minimum Energy Requirements for Space Travel¹

By

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(With 12 Figures)

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Abstract - Zusammenfassung - Résumé

Minimum Energy Requirements for Space Travel. The minimum energy requirement for many space missions is calculated and expressed as velocity requirement of a rocket vehicle, which is supposed to fulfill them. This gives a fast possibility for a preliminary outline of an optimum vehicle, or for an approximation of the payload capability of a given vehicle.

Of course, mission flight times can be reduced by utilizing more than minimum energy. This is particularly pronounced in lunar and interplanetary transfers.

Minimalenergie-Erfordernisse beim Raumflug. Die Minimalenergie-Erfordernisse für viele Raumflugprojekte werden berechnet und als Geschwindigkeitserfordernisse eines Raketenfahrzeuges ausgedrückt, das diese Bedingungen erfüllen soll. Dies bietet die rasche Möglichkeit für einen vorläufigen Entwurf eines optimalen Fahrzeuges oder für die angenäherte Berechnung der Nutzlast-Aufnahm fähigkeit eines gegebenen Fahrzeuges.

Selbstverständlich können die einem Projekt zugeordneten Flugzeiten verringert werden, wenn mehr als die Minimalenergie verwendet wird. Dies wird besonders deutlich bei Übergangsbahnen zum Mond oder zwischen Planeten.

Spécifications minimum d'énergie pour voyages interplanétaires. Les spécifications minimum d'énergie sont converties en spécifications de vitesse de la fusée destinée à remplir une mission. Ceci permet une estimation rapide de la charge payante et de la configuration optimale.

Il est évident que les temps de vol peuvent être réduits par un excédent d'énergie, particulièrement pour les vols lunaires et les transferts interplanétaires.

¹ Statements and opinions are to be understood as individual expressions of the author and do not necessarily reflect the views and opinions of ABMA.

This study is concerned with some of the physical fundamentals of space flight. It is in no way related to any project now being worked on by the Army, nor should it be construed as a description of any future project which may be assigned to the Army.

There is no reference made to any specific hardware now in the making.

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I. Introduction

For an idealized step rocket moving along a straight line free of exterior forces holds the well-known equation

$$V_{id} = \sum_{i=1}^{n} c_i \ln r.$$

This velocity shall be called "ideal velocity capability of the vehicle."

A real rocket moving near the Earth's surface does not reach Vu as there are velocity losses due to drag and the gravitational field. It is easy to show that despite the increase in altitude, which gives a potential energy gain, there is still an energy loss due to gravity. The magnitude of losses depends upon the geometry of the ascent path and on the flight program. In space-flight missions involving high acceleration propulsion systems (order of one g)—and only such shall be of interest here—the usual flight method is to impart to the vehicle a certain mission dependent energy E_1 (say, escape energy), make a correction on the way to the target, and apply the terminal maneuver at the target (perhaps brake the energy E_2 of impact on the moon). The energy E_1 corresponds now to a velocity V_1 and if we add some empirical or semi-empirical corrective figures for drag and gravity loss, usually taking account of the variation of specific impulse with ambient pressure by taking some convenient mean value for jet velocity in the first stage, we can quote an ideal velocity figure which is necessary for the vehicle to have in order to perform the escape: "ideal velocity requirement of the mission." Adding to this value for escape some empirical number for correction, and a figure representing E_2 , we get the ideal velocity requirement for a soft lunar landing.

If a vehicle shall do a certain mission, then, of course, the ideal velocity capability of the vehicle must be larger, or equal to, the ideal velocity requirement of this mission.

In those requirements a correction is necessary since we can get from Earth rotation up to 460 m/sec free; this means that the "ideal requirement" can be 460 m/sec less than one would think from the pure energetic point of view. But often because of flight geometry and other considerations only 300 m/sec can be utilized.

In order to have a yard stick for the vehicle capabilities, where are we now? As known, lunar space probes have been successfully fired, meaning that $V_{id}=12.45~\rm km/sec$ is available. As there is much talk about a Venusian space probe, $V_{id}=13.0~\rm km/sec$ can be expected. What is the ultimate limit of chemical rocket vehicles? Assuming a launch weight of $10\cdot 10^6~\rm lb$, and a gross payload (including guidance, etc.) of $1000~\rm lb$, we have a growth factor of 10^4 . For a four-stage vehicle this means a payload ratio per step of 0.1. Assuming a structure ratio of 0.04, we have for one stage a mass ratio of about 7.14, leading with a jet velocity of $4.5~\rm km/sec$ to $V_{id}=35.4~\rm km/sec$. Applying such tricks of the trade as orbital technique, the figure may rise to $45~\rm km/sec$, with some small increases possible due to the use of orbital technique at the target. As perturbation maneuvers can reduce the requirements slightly, we can understand them to be a small increase in velocity capability. The "feasibility limit", then, can somewhat arbitrarily be set perhaps at $V_{id}=50~\rm km/sec$.

As the problems of such a vehicle are stupendous (as all of you realize), I feel, therefore, that a practical limit will be lower. Fortunately, we do not need such high ideal velocities within the solar system, as the "principle of mission staging" is a way to avoid too high requirements. For example, you go to the

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lize), I ot need nission to the lunar surface using many landing vehicles, and assemble there an Earth-return vehicle out of the useful payload of your landers. So you can return to Earth, using several vehicles of one-way capability instead of one vehicle of two-way capability. Therefore, as an estimate, the practical limit for a chemically-propelled rocket vehicle may be near $V_{ii}\!=\!25~{\rm km/sec}$.

Using a "conventional" nuclear propulsion system as it is seen today (nuclear reactor heating a working fluid, which expands through a nozzle), we could perhaps get 80 km/sec instead of the 50 for the chemical system. The more practical limit could conceivably be near 30 km/sec.

The low-acceleration systems could have still higher velocity capabilities, but part of this is used up against the higher gravity losses. Perhaps as ball-park figures, 100 km/sec as ultimate limit, and 50 km/sec as practical limit, can be envisioned.

To meet still higher requirements, very exotic propulsion systems must be used; for example, Prof. Sänger's photon_propulsion.

II. Velocity Requirements for Earth-Bound Missions

For comparision only, some approximate figures have been computed, which give range versus ideal velocity capability for ballistic-type missiles (Table I).

III. Velocity Requirements for Satellite Missions (Circular Orbits)

A. Nonrotating Earth (Polar Orbits)

Assuming a Hohmann transfer (Fig. 1), a first kick is necessary to throw the payload up to orbital altitude:

$$\begin{split} \delta_{1}{'} &= \left(\frac{\gamma\,M}{R}\right)^{\frac{1}{2}} \left(\frac{2\,r}{r+R}\right)^{\frac{1}{2}}. \end{split}$$
 Introducing $\frac{\delta'}{\left(\frac{\gamma\,M}{R}\right)^{\frac{1}{2}}} = \delta$, and $\frac{r-R}{R} = k$ comes

$$\delta_1 = \left(\frac{2+2k}{2+k}\right)^{\frac{1}{2}}.\tag{1}$$

The arrival velocity at apogee is

$$V_{\bullet} = \left\{ \frac{2}{(2+k) \ (1+k)} \right\}^{\frac{1}{2}}.$$

So a second kick to circularize must be given

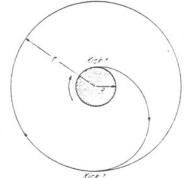


Fig. 1. Hohmann transfer. Circular velocity for E=R equals $[(\gamma M)/R]^{\frac{1}{2}}=7910$ m/sec for Earth

$$\delta_2 = \frac{1}{(1+k)^{\frac{1}{2}}} \left\{ 1 - \left(\frac{2}{2+k} \right)^{\frac{1}{2}} \right\}. \tag{3}$$

The total ideal requirement for the mission then is

$$\Delta_1 = \delta_1 + \delta_2 = \frac{k + (1 + k/2)^{\frac{1}{2}}}{(1 + k)^{\frac{1}{2}} (1 + k/2)^{\frac{1}{2}}} \approx 1 + k/2, \text{ for small } k.$$
 (4)

Introducing for the moment $(1 + k/2)^{\frac{1}{2}} = x$, this can be written

The solution is $x \approx 2.88$, leading to k = 14.58176 for the most difficult circular orbit with $\Delta_1 \approx 1.5362$.

Using a three-kick transfer (Fig. 2), the first kick is used to throw the payload to infinity.

$$\delta_1 = \sqrt{2} \ . \tag{5}$$

The second kick at infinity is a zero-adjustment kick; the third kick brakes the arrival speed (escape) to circular speed:

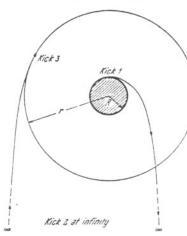


Fig. 2. Three-kick transfer

$$\delta_3 = (\sqrt{2} - 1) \frac{1}{(1+k)^{\frac{1}{2}}}.$$
 (6)

In total, we have used

$$\Delta_2 = \sqrt{2} + \frac{\sqrt{2} - 1}{(1+k)^{\frac{1}{2}}}.$$
 (7)

Comparing Δ_2 and Δ_1 , there is $\Delta_2 < \Delta_1$ for k > 10.94, which means, that the three-kick transfer is superior for very high orbits.

So, energy-wise, the circular orbit at k=10.95 is the most difficult one, with $\Delta=1.5340$, if the three-kick transfer is used for larger altitudes. (These considerations are of theoretical value only, since k=10.94 corresponds to satellite

altitudes of no interest, and besides the three-kick transfer has other serious disadvantages: Long travel time, and extreme sensitive to the second kick.)

B. Rotating Earth

If the firing occurs from a latitude λ under an azimuth a (East of North) then there is an assistance given from Earth rotation of approximately

$$V_{g} = 450 \cdot \cos \lambda \sin \alpha \, [\text{m/sec}].$$
 (8)

C. Change of Orbital Plane

Only a simple case shall be considered here; viz, to go into an equatorial orbit by launching from the latitude φ under an azimuth of $a=90^\circ$. Ascent is via Hohmann transfer to a waiting orbit at 200 km altitude. This is necessary in order to place the apogee of the transfer ellipse from the waiting to the final orbit over the equator.

$$\Delta_1 = 1.0157 \approx 1 + \frac{k}{2}$$
.

An escape from there would require $\Delta_{\rm Esc}=(\sqrt{2}-1)\left(1-\frac{k}{2}\right)$. So the total escape $\Delta=\sqrt{2}+(2-\sqrt{2})\,\frac{k}{2}$.

The direct escape is $\sqrt{2}$ —so there is a waiting-orbit loss for escape of

$$\varDelta_{\scriptscriptstyle L} = (2 - \sqrt{2}\,)\,\frac{k}{2} \approx 0.0092$$
 .

Because of maneuvering, I will take $\Delta_L = 0.01$ as "typical waiting-orbit loss."

The upper kick for circularization is given by

 $\delta = \frac{\sqrt{1 + \frac{k}{2}} - 1}{\sqrt{1 + \frac{k}{2}} \sqrt{1 + k}} \,. \tag{9}$

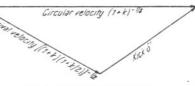


Fig. 3. Change of inclination

If there is a plane change φ involved (see Fig. 3), then the kick to both circularize and accomplish the plane change is

$$\delta^{1} = \sqrt{\delta^{2} + \frac{4\sin^{2}\varphi/2}{(1+k)\sqrt{1+k/2}}} . \tag{10}$$

The total velocity required on a rotating Earth then is

$$\Delta_{1\varphi} = \Delta_1 - \frac{450}{\left(\frac{\gamma M}{R}\right)^{1/2}} + (\delta^1 - \delta) + \frac{900}{\left(\frac{\gamma M}{R}\right)^{1/2}} \sin^2 \frac{\varphi}{2} + 0.01 S(\varphi)$$
 (11)

where $S(\varphi) = 0$ for $\varphi = 0$ and

 $S(\varphi) = 1$ for $\varphi = 0$

 $[\delta', \delta \text{ are given by eqs. (10, 11); } \Delta, \text{ by eq. (4)}].$

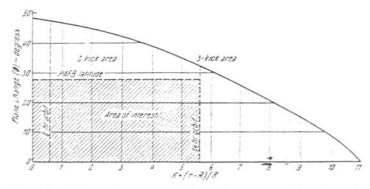


Fig. 4. Three-kick transfer versus two-kick transfer (energy considerations only)

The corresponding expression for the three-kick transfer is

$$\Delta_{2\varphi} = \Delta_2 - \frac{450}{\left(\frac{\gamma \dot{M}}{R}\right)^{1/2}} + \frac{900}{\left(\frac{\gamma \dot{M}}{R}\right)^{1/2}} \sin^2 \frac{\varphi}{2}$$
 (12)

[Δ_2 is given by eq. (7)].

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t is via ssary in he final Energywise, we consider the function

$$A\left(\varphi,k\right) = \Delta_{2\varphi} - \Delta_{1\varphi}. \tag{13}$$

If A > 0, then the Hohmann transfer is preferable; A < 0, then the three-kick transfer is preferable.

It is easy to show:

1. $k < 10.94, \varphi = 0: A > 0$

2. $k = \infty$, φ arbitrary: A = 0

3. $k = \text{arbitrary}, \, \varphi > 48.3^{\circ} : A < 0$

For $A(\varphi, k) = 0$, (see Fig. 4).

It is seen that the three-kick transfer appears not to have a field of practical application.

D. Examples

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Circular 568 km-orbit, 96-min, equatorial: Equator-launched, Hohmann Equator-launched, three-kick $\varphi = 28^{\circ}\text{-launched, Hohmann}$ $\varphi = 28^{\circ}\text{-launched, three-kick}$	$\begin{array}{l} (\gamma~M/R)^{1/2} = 7.920~\text{m/sec} \\ V_{id} = 0.99 = 7.841~\text{m/sec} \\ 1.76 = 13.939~\text{m/sec} \\ 1.44 = 11.405~\text{m/sec} \\ 1.766 = 13.987~\text{m/sec} \end{array}$
Circular 24-hr orbit, equatorial: Equator-launched, Hohmann Equator-launched, three-kick $\varphi = 28^{\circ}\text{-launched, Hohmann}$ $\varphi = 28^{\circ}\text{-launched, three-kick}$ $\varphi = 28^{\circ}\text{-launched, three-kick}$ $\varphi = 28^{\circ}\text{-launched, 28}^{\circ}\text{ inclined orbit}$	1.45 = 11,484 m/sec 1.52 = 12,038 m/sec 1.51 = 11,959 m/sec 1.526 = 12.086 m/sec 55 m/sec to above equator-launched velocities
Escape: equatorial-launched, Eastward	10,736 m/sec 11.186 m/sec

What does this mean in payload? Let us assume the following arbitrary vehicles for the first mission:

	Type					
	A	В	C			
Payload (Container) Weight	100	80	70			
Guidance, Control, Instruments	4	11	14			
Fuselage	4	11	14			
Motor, etc.	4	10	14			
Total	112	112	112			

We will look at the HOHMANN-transfer missions only

We will look at the HOHMANN-transfer missions only
$$28^{\circ}$$
-launched, 28° inclined orbit: with $I_{sp} = 300$ sec, $55 = 300$ g. $\ln \frac{M}{m} \rightarrow \frac{M}{m} = 1.0186$, which leads to $m = \frac{112}{1.0186} = 109.95$, or fuel used: 2.05 Additional Tankage: 0.2

Payload loss: 2.25 of 100 for Vehicle A.

With an I_{sp} of 400 sec, the payload loss was 1.65 of 100. In this manner the following table has been computed.

Payloads Including Containers
(HOHMANN transfers only)

		$I_{sp} = 300 \text{ sec}$			$I_{sp} = 400 \text{ sec}$	
	A	В	C	A	В	· C
Equator-launched						
Equatorial 96-min Orbit 28°-Launched,	100.0%	100.0%	100.0%	100.0%	100.0%	100.0%
28° Inclined	97.8%	97.0%	96.9%	98.4%	97.9%	97.6%
Equatorial Equator-launched,	13.4%	-	_	26.4%	8.0%	-
Equatorial 24-hr Orbit 28°-Launched, 28° In-	12.4%	-	_	26.0%	7.5%	-
clined 24-hr Orbit 28°-Launched,	11.9%	- •	_	24.9%	6.1%	_
Equatorial 24-hr Orbit Escape, Equator-	7.2%	-	_	19.9%	0.0%	_
launched, Due East Escape 28°-launched,	22.8%	3.5%	_	35.6%	19.5%	8.0%
Due East Escape Polar	22.0 % 16.2 %	2.5%	_	34.9%	18.6% 11.5%	7.0%

This table illustrated several well-known facts:

1. If azimuth of firing is 90° , then only small penalties are involved in going up to 30° off the equator.

2. If higher energy missions are flown, then it is not optimum just to exchange propellant for payload in the last stage. But if this is done, then the result is very dependent upon specific impulse.

 Even for vehicles of similar performance at one mission, the performance at other missions may vary widely.

E. Recovery of Satellites

Here one kick is assumed to brake the orbital velocity so far that a transfer ellipse is entered the pericenter of which is sufficiently deep in the atmosphere, so that further braking is done by atmospheric drag. Final descent could employ lift, or a parachute, or in the case of some types of instruments simply impact.

F. Conversion of "Ideal Required Velocity" to "Required Velocity"

Certain correction figures have to be added, which are of an empirical nature. Those correction figures take care of:

- 1. Gravity loss and drag loss: 1500 m/sec seem to be a usable figure for escape missions and large vehicles. (Of course, this number depends on vehicle shape. trajectory and acceleration program.)
 - 2. Maneuvering Reserve: This has to be estimated for every mission.
- 3. Unusable propellant residuals and mixture ratio shifts: With control of mixture ratio and trapped propellants, 3% of the ideal velocity should suffice.

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IV. Velocity Requirements for Space Probes

By space probe a vehicle is meant which is used for space research without necessarily approaching a planet or any other celestial body. Therefore, the guidance and space navigation problems are greatly simplified. To go from Earth to the Moon, escape velocity $\left(2 \frac{\gamma M}{R}\right)^{1/2}$ is approximately necessary, the minimum for direct transfers being about 100 m/sec lower. If we look for the circular orbit of the same energy requirement, we have to solve $\frac{k+\sqrt{1+k/2}}{\sqrt{(1+k)\,(1+k/2)}}=\sqrt{2}$, from which comes $k \approx 2.303$.

So we can conclude, that, for a simple Earth-escape experiment, about the same payload can be carried as into the circular orbit of k = 2.303. In practice even more can be carried, as the guidance system should be simpler and, therefore,

For interplanetary probes, obviously, more energy is required than for probes in near-Earth Space. It is easy to show, that, disregarding Earth, a minimum perihelion velocity of 32.83 km/sec at Earth distance (149.106 km) from the Sun is necessary in order to place the aphelion out to 230.106 km (Mars distance).

This is 32.83 - 29.8 = 3.03 km/sec above local circular velocity (Earth velocity, going around Sun). Therefore, the minimum required launch velocity is

 $\sqrt{11.1862^2 + 3.03^2} = 11.5893 \text{ km/sec}$

which is only 403.1 m/sec more than

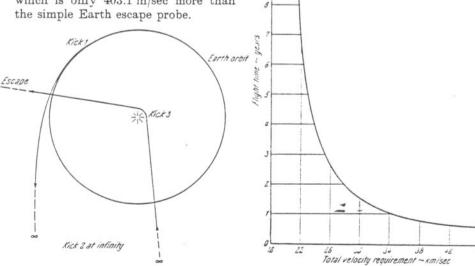


Fig. 5. Solar system escape - Hi-performance vehicle

Fig. 6. Solar probe: Flight time versus total

For Venusian probe, put 2.4 km/sec instead of the 3.03 km/sec for Mars, resulting in $V_{min} = 11.4408 \text{ km/sec}$, only 254.6 km/sec above escape velocity.

The minimum requirement for a Mercury probe is 13.504 km/sec or 2.318 km/sec excess over Earth escape.

The minimum to escape the solar system is to have a residual velocity of $(\sqrt{2}-1)\cdot 29.8=12.3436$ km/sec. This leads to a total minimum requirement of 16.6582 km/sec which is 5.472 km/sec over simple Earth escape.

Assuming we had a vehicle of a capability of $V_{id} = 50$ (100) km/sec. To leave the solar system (Fig. 5) it would be best to apply a first impulse in order just to

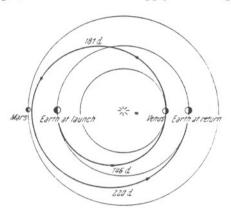


Fig. 7. Hohmann's round-trip trajectory

M	aneuvers:		
1.	Brake original 29.8 km/sec of Earth by	2.4 km/se	c
2.	Increase on passing Venus to 39.4 km/sec by	1.8 km/se	ю
	Increase on passing Mars to 24.8 km/sec by		
4.	Brake 31.5 km/sec arrival at Earth by	1.8 km/se	c
	Total velocity requirement	17.6 km/se	e
	Total flight time	547 days	

leave the system: $16.66 + losses \approx 20$ km/sec. At infinity, zero kicks would adjust for a return to just graze the Sun. Arrival velocity would be local escape velocity = =617.5 km/sec at 700,000 km from the Sun's center. To this we add the remaining 30 (70) km/sec, leading to a velocity remaining at infinity of 195 (302) km/sec. Travel time to the nearest fixed star, about 4 light-years distant, would now be about 6150 (3980) years, or total mission time about 6200 (4000) years. Therefore, with the highly advanced vehicles which were assumed, not even the nearest fixed stars are within reach. (In order to get to them in 40 years of transit time, the ideal velocity capability of the vehicle must be of the order of 30,000 km/sec.)

In solar probes we have to differentiate between three classes:

1. Probes in near solar space. A good example of this is the Mercury probe, for which we found $V_{\iota\iota}$ min = 13.504 km/sec. The flight time is about 111 days.

2. Direct probe to the Sun: The Earth orbital velocity is braked, and the vehicle drops in vertically to the Sun. The required ideal velocity is high—31.8 km/sec with a flight time of only 65 days.

3. Indirect solar probe. Here a solar system escape must be performed first and then a drop towards the Sun. The ideal energy required is only 16.66 km/sec, but the flight time is infinite. For practical cases, you would, of course, not go quite to infinity, resulting in a higher velocity and lower flight time required (see Fig. 6).

More sophisticated types of probes are those performing round trips, two of which may be of special interest:

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 HOHMANN Roundtrip (Fig. 7). CROCCO Roundtrip (Fig. 8).

As these probes can conceivably return to Earth, a manned version of such an expedition may be interesting.

V. Lunar Flights

Six types of flights must be considered:

- 1. Probes
- 2. Hard impacts
- 3. Circumlunar Flights
- 4. Lunar Satellites
- 5. Soft Landers
- 6. Earth-Moon Return Flights.

Probes and hard impacts differ mainly in guidance accuracy required. Therefore, no more additional information seems necessary here.

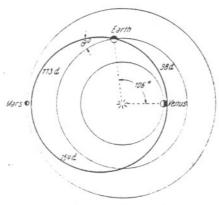


Fig. 8. Crocco's round-trip trajectory

Total velocity requirement	15.968 km/sec
Total flight time	12.5 to 13.5 months
Possible launch time	June 1971
J	16.040

For the circumlunar flight the energy requirement, apart from maneuvering, is similar to Earth escape. If a return to Earth is planned, then, because of the critical atmospheric re-entry, ample maneuvering fuel should be provided for.

Let us look for a moment at some characteristic data for transfer trajectories in the Earth-Moon system:

Injection angle: near horizontal

	Cutoff Velocity in Inertial Earth-Centered Space at 200 km Altitude	Flight Time	Unbraked Luna Impact Velocity		
1 .	10,881 m/sec	~10 years	2,325 m/sec		
2	10,920 m/sec	~ 5 days	2,500 m/sec		
3	10,970 m/sec	21/2 days	2,705 m/sec		
4	11,015 m/sec	51 hrs	2,886 m/sec		
5	11,100 m/sec	41 hrs	3,179 m/sec		

Remarks:

1. Absolute minimum injection velocity to reach the Moon, from Jacobi's Integral.

2. Minimum for direct Earth-Moon transfer.

3. Two-and-a-half day transfer.

4. Injection velocity equal local escape velocity.

5. "Fast" trajectory.

Some Data:

Some Data.		
Mean Earth Radius	6,371.1 km	
Mean Lunar Radius	1,738 km	
Escape velocity, Earth, zero altitude	11,186 km/s	ec
Escape velocity, Moon	2,374 km/s	ec
"Mean" Earth—Moon distance		

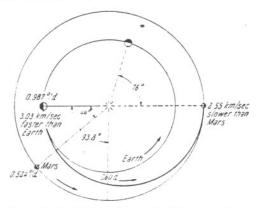


Fig. 9. Earth—Mars mission profile (Hohmann ellipse) Opposition occurs (44)/(0.987-0.524)=95 days $=93.8^{\circ}$ after launch. The waiting time is in this case (360-2.76)/(0.987-0.524)=449 days. So a return mission last (260/449/260)=970 days

From the guidance point of view, the "slow" trajectories show large deviations for small injection errors. Therefore, trajectory No. 4 seems to be a good compromise, with 3 being a competitor.

Upon arrival for a lunar satellite, we have to brake:

Satellite Altitude	Trajectory	Circular Velocity	To be Braked
$\frac{\dot{h}}{R_{Moon}} = 0$	$2^{1}/_{2}$ days	1678 m/sec	1027 m/sec
0	51 hrs	1678 m/sec	1208 m/sec
$0.25 \\ 0.25$	$2^{1}/_{2}$ days 51 hrs	1502 m/sec 1502 m/sec	985 m/sec 1181 m/sec
1	21/2 days	1188 m/sec	932 m/sec
1	51 hrs	1188 m/sec	1160 m/sec

A soft lunar landing vehicle has to brake the total impact velocity by rocket action.

The return flight is, energy-wise, much simpler; because, at Earth's side probably the atmosphere can be used for re-entry braking. However, corrective fuel should be provided as the ideal re-entry conditions must be met rather closely.

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VI. Orbital Technique

It is necessary to say a few words on the use of orbital technique.

a) Upon Departure

For example, in the 96 min-orbit a space vehicle could be assembled. This, then, could take advantage of the energy it already has, and a noticeable structural advantage should result from the possibility of using relatively low accelerations, and from the absence of aerodynamic considerations. Maneuvers into the 96-min

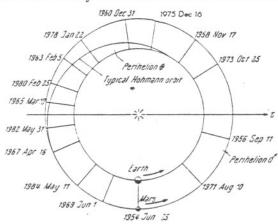


Fig. 10. Opposition of Mars and HOHMANN ellipse injection dates

		T. 1	5. 10.	Opposition	Or Trong o	114 11	OHJIA	empse mject	non c	acca
HOHMANN	ellipse	inj	ection	dates:						
		23	Aug	1958	23	Feb	1969	7	Oct	1977
		25	Sep	1960	21	May	1971	29	Oct	1979
		16	Oct	1962	7	Aug	1973	4	Dec	1981
		9	Nov	1964	12	Sep	1975	28	Jan	1984
		25	Dec	1966						

Notes: Maximum error in injection dates: 1958—1971, 3 days; 1971—1984, 5 days. The opposition dates were obtained by extending the ephemeris of the Earth and Mars by a numerical integration method. This ephemeris was extended through 1984 by use of a graph showing synodic period of Mars versus longitude of opposition, prepared from the American Ephemeris and Nautical Almanac. The injection dates were obtained by use of a graph showing opposition time minus injection time versus longitude of opposition. This graph was prepared by choosing a multiple of conjunction times of the HOBMANN ellipse and the orbit of Mars and from this, calculating the period and the theoretical opposition times and longitudes. These were then plotted to obtain the graph

orbit and the resulting high cutoff altitude for the powered phase leaving the orbit will result in some small energy losses.

b) Upon Arrival

At the target site, going first to an orbit and from there down to the surface has potential advantages:

- 1. Perhaps better control of landing area.
- 2. Higher safety of mission success, even if the landing fails.
- 3. Saving of energy, if return fuel is left in orbit and picked up during the return phase.

Disadvantages are:

- 1. More complicated over-all scheme.
- 2. All braking could be done aerodynamically if there is an atmosphere; if the braking to orbit is done by rocket, usually much more unfavorable conditions exist.

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sembled. This, able structural accelerations, ato the 96-min

. The opposition rical integration ic period of Mars al Almanac. The stion time versus tion times of the theoretical opporaph

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o the surface

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mosphere; if le conditions

c) Upon Return

Upon return to Earth, the following possibilities exist:

1. Direct re-entry to the atmosphere.

2. Return to 96-min orbit, thereby using rocket braking. The disadvantage of using fuel is partly compensated for by the advantage of using a special vehicle for transferring and receiving.

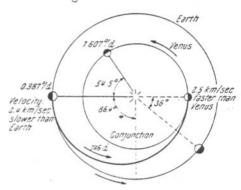


Fig. 11. Earth—Venus mission profile (HOHMANN ellipse)

Inferior conjunction is (54.5)/(1.607-0.987)=87.6 days after launch; the angle is $87.6\cdot0.987=86.4^\circ$. From this follows the rule: Launch $86.4^\circ=87.6$ days before inferior conjunction. For the return flight, we need to have Earth leading by 36° . (In above picture, we simply reverse all directions for the return flight.) So Venus has to catch up $360-2.36=298^\circ$, which takes $(288)/(1.607-0.987)\approx 460$ days. So a return mission lasts 146/480/146=750 days

3. Return to another orbit, being picked up there and brought to the 96-min orbit, from where transportation to the Earth's surface is provided. A very interesting return orbit of this type is the following elliptic orbit, which often needs only little rocket braking:

Elliptic Orbit Properties

 $Perigee: 568 \; km \; altitude \; (circular \; velocity: 7581 \; m/sec, \; escape \; velocity: 10,721 \; m/sec).$

Velocity: 10,431 m/sec (2850 m/sec above circular velocity).

Apogee: 120,143.4 km from Earth's center.

Eccentricity: 0.89321044. Major axis: 65,044.7 km. Minor axis: 29,247.0 km. Period: 45.848 hr.

So all the braking necessary is 10.721-10.431=290 m/sec. Of course, now the pick-up vehicle has to have an ideal minimum capability of $2\times2850=5700$ m/sec.

VII. Planetary Flights

Corresponding to paragraph V, six mission types exist. Only planetary satellites, soft landers and return flights are of interest, as probes are already treated elsewhere.

Perhaps advanced propulsion systems—e.g. ionic systems—will be used for large-scale interplanetary operations, but we will limit ourselves here to the conventional chemical impulsive systems.

Only four representative missions via Hohmann transfers are considered in some detail, bearing in mind that the return flight has somewhat symmetrical demands.

a) Martian Satellite

As shown, the minimum ideal initial (Earth-side) launch velocity equals 11.5893 km/sec. Mars has to be there, when the vehicle reaches the Aphel of its orbit. The transfer time from Earth is 260 days. Mars moves during these 260 days through $0.524 \cdot 260 = 136$ degrees of arc. Therefore, at launch Mars must be 44 degrees ahead of Earth (which is $\frac{44}{0.987-0.524} = 96$ days before opposition). -The Martian orbit is inclined by 1°51' to the ecliptic. If the vehicle shall go into the Martian plane of motion at the node, the ideal velocity requirement is about 27 · sin 1°51′ = 0.872 km/sec. By giving the plane change kick about halfway and arriving at Mars moving not within the Martian plane, usually (if the node does

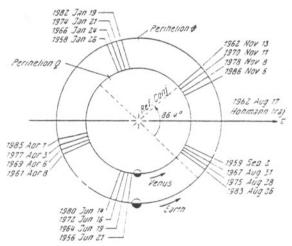


Fig. 12. Inferior conjunction of Venus and Hohmann ellipse injection date

Hohmann ellipse injection dates:

,								
1956 M	ar 26		1967	Jun	4	1978	Aug	13
1957 Oc	ct 31		1969	Jan	8	1980	Mar	18
1959 Ju	ın 6		1970	Aug	15	1981	Oct	23
1961 Ja	n 10		1972	Mar	21	1983	May	30
1962 At	ug 17	19	1973	Oct	26	1985	Jan	3
1964 M	ar 23		1975	Jun	2	1986	Aug	10
1965 Oc	t 28		1977	Jan	6	_	-	

Notes: 1. Estimated error of a conjunction or a Hohmann injection date is 4 days maximum.

2. Orbits and circles; $r(\oplus) = 1.00$ au, r(?) = 0.723 au.
3. Dates and longitudes of inferior conjunctions are: $t = \text{JD } 243\ 5062.3/583.931\ \text{N (days)}$; and $t = 575^\circ$ 5180 N; Reference date = 1954 Nov 15.29 UT.

4. Constant synodic period = 583.921 (days). 5. Hohmann date and longitude; t=87.6 days, $\lambda=86^{\circ}.4$

not happen to be at this place) some savings can be accomplished. If the vehicle travels neither in the ecliptic nor in the Martian plane, and if the central angle between launch and arrival is slightly smaller than 180 degrees, then the loss due to inclination change becomes negligible. (Oral communication from Dr. D. F. LAWDEN.)

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Arrival at the Martian orbit occurs with a velocity of 2.55 km/sec less than Mars orbital velocity. At 1000 km above the Martian surface, the vehicle will move 5.15 km/sec; an ideal minimum braking of 2.01 km/sec brings this down to 3.14 km/sec, which is local circular velocity (escape velocity at zero altitude: 5.04 km/sec).

b) Venusian Satellite

At launch, Venus must be 54.5 degrees behind Earth (or 88 days before inferior conjunction). After 146 days, the vehicle approaches Venus, being 2.5 km/sec faster than Venus. As the inclination of the Venusian orbit is about $3^{\circ}24'$, the maximum requirement of the ideal velocity for the change is $32.5 \sin 3^{\circ}24' = 1.93 \text{ km/sec}$.

Escape velocity at Venusian surface is 10.23 km/sec; at 1000 km altitude, this is reduced to 9.49 km/sec (circular velocity: 6.71 km/sec). The vehicle arrives with $\sqrt{(2.5)^2 + (9.49)^2} = 9.814$ km/sec. Therefore, a braking of 3.104 km/sec is necessary.

c) and d) Soft Landings on Mars or Venus

Here I will always assume, that the atmosphere is used for braking and landing. The speed which has to be broken upon arrival is about

Venus: 10.5 km/sec (from Earth) Earth: 11.5 km/sec (from Mars, Venus) Mars: 5.64 km/sec (from Earth)

If this is done aerodynamically, then there is only a little fuel used for control. (For the timing and outlay of Martian and Venusian flights, see Figs. 9—12.)

Of some interest might be a manned planetoid mission, because this would give the chance of actually being on another star and doing research there. For a planetoid within Mars' orbit and of negligible gravity, appropriate Mars data can be used.

Appendix

Table I. Ideal Velocity Requirement for Earth—Bound Missions of Ballistic—Missile Types

Range (km)	V_{id} (km/sec)
500	3
1,000	3.9
2,000	5
5,000	6.8
10,000	8.5
20,000	8.8

Remarks: According to paragraph III—F, 3% of V_{id} may be added for mixture ratio shifts, trapped residuals and flight performance reserves.

n. = 575°

hicle angle due from

Table II. Velocity Requirements for Some Missions (See Remarks on Table I)

	Equator	al Earth	Satellite) al				· Lunar Missions							
		Equator	Launch		Launch		Launch	1000	Satellite		Circum- Lunar	+ Return Satellite	Soft Landing		
	Values in km/sec	96-min	24-hr	96-min	24-hr	96-min	24-hr								
ch	Ideal Minimum Launch	8.291	11.934	11.801	12.359	8.291	11.934	11.186	11.186	11.186	11.186	11.186	11.186		
launch	Rotational Gain	0.45	0.45	0.4	0.4	0.45	0.45	0.3	0.3	0.3	0.3	0.3	0.3		
7	(7-Loss	1.4	1.42	1.4	1.42	1.4	1.42	1.42	1.42	1.42	1.42	1.42	1.42		
Earth	Drag-Loss	0.15	0.16	0.15	0.16	0.15	0.16	0.16	0.16	0.16	0.16	0.16	0.16		
ñ	Maneuvering	0.05	0.05	0.01	0.01	0.05	0.05	0.05	0.05	0.05	0.05	0.05,	0.05		
	Maneuvering Transfer		0.05		0.05	******	0.05	0.05	0.05	0.05	0.05	0.05	0.05		
7									•						
riv	Ideal Minimum			*					1.18	2.89		1.18	2.89		
8.1	Rotational Gain	-				N						process			
Pee	<i>G</i> -Loss						N		0.03	0.3		0.03	0.3		
At target arrival	Maneuvering	0.05	0.05	0.05	0.05	0.05	0.05		0.05	0.05	0.1	0.05	0.05		
-4															
ich	Ideal Minimum Launch					0.06	1.49			W- M- 10		1.18	2.89		
launch	Rotational Gain		11.1904		10.00							*****	0.2		
1,	<i>G</i> -Loss											0.03	0.15		
Target	Drag-Loss	*** *	p	900 - 100	0.00		\$170.0FM			M-100000		******			
Ta	Maneuvering		P					_			_	0.05	0.05		
	Maneuvering Transfer										0.05	0.05	0.05		
1	Earth Landing Manouver					0.05	0.05				0.05	0.05	0.05		
İ	Total	9.50	13.20	13.10	13.75	9.60	14.75	12.55	13.85	15.80	12.75	15.20	19.35		

						1	-		-		0.00	0.05	0.03
i	Earth Landing Maneuver	30.00	13.20	13.10	13.75	0.05 9.60	$0.05 \\ 14.75$	12.55	13.85	15.80	0.05 12.75	$0.05 \\ 15.20$	0.05 19.35

Table II (Continued)

				Space Pr	obes				
			No Re	ecovery		HOHMANN	Crocco	Simple Solar	Solar
	Values in km/sec	Simple Earth Escape	Martian Probe	Venusian Probe	Mercury Probe	Probe + Return	Probo + Return	System Escape	Probe
Earth launch	Ideal Minimum Launch Rotational Gain (I-Loss Drag-Loss Maneuvering	11.2 0.4 1.42 0.16 0.05	11.589 0.3 1.42 0.16 0.10	$11.441 \\ 0.3 \\ 1.42 \\ 0.16 \\ 0.10$	13.504 0.3 1.45 0.17 0.10	11.441 0.3 1.42 0.16 0.10	13.948 0.3 1.45 0.17 0.10	16.658 0.4 1.46 0.18 0.10	17.0—31.8 0.4 1.46 0.18 0.10
	Maneuvering Transfer								_
At target arrival	Ideal Minimum Rotational Gain G-Loss Maneuvering		0.2	0.2	0.2	0.6 + 1.8 + 2.3	0.6	-	0.2
larget launch	Ideal Minimum Launch Rotational Gain G-Loss Drag-Loss Maneuvering	=						=======================================	= ,
	Maneuvering Transfer Earth Landing Maneuver Total	 12.45	13.15	13.0	15.1	0.1 17.60	0.1 16.05	 18.0	 18.55—33.

Minimum Energy Requirements for Space Trave

H. O. RUPPE:

				Planetar	y Missions					
			No Re	ecovery				+ Return		
	Vulues in km/sec	Mars Satellito	Vonus Satellito	Mars Soft Landing	Venus Soft Landing	Mars Satellite	Venus Satellite	Mars Soft Landing	Venus Soft Landing	Planetoid in Mars Orbit, Soft Landing
Earth launch	Ideal Minimum Launch Rotational Gain G-Loss Drag-Loss Maneuvering	$11.589 \\ 0.3 \\ 1.42 \\ 0.16 \\ 0.05$	$11.441 \\ 0.3 \\ 1.42 \\ 0.16 \\ 0.05$	$11.589 \\ 0.3 \\ 1.42 \\ 0.16 \\ 0.05$	11.441 0.3 1.42 0.16 0.05	11.589 • 0.3 1.42 0.16 0.05	$11.441 \\ 0.3 \\ 1.42 \\ 0.16 \\ 0.05$	11.589 0.3 1.42 0.16 0.05	11.441 0.3 1.42 0.16 0.05	11.589 0.3 1.42 0.16 0.05
	Maneuvering Transfer	0.4	0.5	0.4	0.5	0.4	0.5	0.4	0.5	0.4
At target arrival	Ideal Minimum Rotational Gain G-Loss Maneuvering	2.01 0.1 0.05	3.104 	0.01	0.01	2.01 	3.104 	0.01	0.01	2.55
larget launch	Ideal Minimum Launch Rotational Gain G-Loss Drag-Loss Maneuvering					2.01 0.05 0.05	3.104 0.1 0.05	5.64 0.2 0.3 0.15 0.05	10.5 1.5 0.2 0.05	2.55 — — 0.05
	Maneuvering Transfer Earth Landing Maneuver Total	15.50	16.60	13.40	13.35	0.04 0.05 18.05	$0.05 \\ 0.05 \\ 20.40$	0.04 0.05 19.80	0.05 0.05 26.15	0.04 0.05 18.95

Table III. Velocity Requirements for Various Missions (Equator-Launched, Unless Otherwise Specified)

On this table the total velocity requirements for a number of missions are summarized. "Direct Target" means, that the maneuver at target is employed without going to an orbit around the target first. "Orbit at Target" on the other hand implies the use of orbital technique at target. The "Elliptical Orbit" is the same as is described in paragraph VI of this report. Also, see remarks on Table I.

		Eq	uatorial	Earth	Satellit	es, Circu	lar	Ellip Equal Earth S	torial	Lunar Missions							
		hod		+	Rocovoi	·y	k k	ret	3	ing	+ Return						
	150 km 1	6 min	24 hr	96 min	24 hr	150 km	96 min	24 hr	150— 35871	150— 35871	Impact	Satellite	Soft Landing	Circum- Lunar		Soft Landing	
Direct Target																	
Earth	9.20	9.50	13.20	13.10	13.75				11.8		12.55	13.85	15.80				
Orbit			3.85		4.40		-				3.10	4.40	6.30				
Earth—Earth						9.25	9.60	14.75		11.9			_	12.75	15.20	19.33	
Orbit—Earth			-					5.4						3.25	5.75	9.70	
Earth—Orbit		Marian.						17.08	-	-				15.90	18.40	22.50	
Orbit—Orbit		-					-	7.73				-		6.40	8.85	12.88	
Earth—Elliptical Orbit		-					*******			-				13.05	15.45	19.50	
Orbit—Elliptical Orbit									_	-				3.55	6.00	10.00	
Orbit at Target																	
Earth		-	a-residua.					**********	_		13.7		15.95				
Orbit											4.25		6.45				
EarthEarth								**********								19.4	
Orbit—Earth																9.8	
Earth-Orbit										*****			-	_	- ,	22.6	
Orbit—Orbit			-		-											12.9	
Earth—Elliptical Orbit																19.6	
Orbit—Elliptical Orbit							***									10.1	

Minimum Energy Requirements for Space Travel

Table III (Continued)

				Spac	e Probe	8			Planetary Missions										
							Ret	urn					+ Return						
	Escape	Martian	Venusian	Mercury	Solar	Solar System Escape	HOHMANN	Скоссо	Mars Satellite	Venus Satellite	Mars Soft Landing	Venus Soft Landing	Mars Satellite	Venus Satellite	Mars Soft Landing	Venus Soft Landing	Planetoid Soft Landing		
Direct Target					18.55											,			
Earth	12.45	13.15	13.00	15.1	33.35 9.20-		_		15.50	16.60	13.40	13.35	_	-		-	15.81		
Orbit	3.10	3.70	3.55	5.65	24.0	8.65			6.05	7.15	3.95	3.90					6.451		
Earth—Earth		-					17.60	16.05					18.05	20.40	19.80	26.15	18.95		
Orbit—Earth							8.15	6.60					8.60	10.95	10.35	16.70	9.50		
Earth—Orbit	-			Program.			21.00	21.95							23.43		22.50		
Orbit—Orbit			******		Normal		11.55	12.50	_	-			12.11	14.31	13.88	20.25	13.05		
Earth-Elliptical Orbit .				***			18.15	19.1					18.71	20.91	20.48	26.85	19.65		
Orbit-Elliptical Orbit .		***************************************			******	Non-service.	8.70	9.65					9.26	11.46	11.03	17.40	10.20		
Orbit at Target																			
Earth				Meta-respo	Marin or	Participal		-	_		15.65	16.75	Manager	-			16.001		
Orbit					-	***			_		6.20	7.30					6.55^{1}		
Earth-Earth															22.00	29.50	19.10		
Orbit—Earth															12.60	20.10	9.60		
Earth—Orbit		-	-						_		*****	-			25.55	32.90	22.60		
Orbit—Orbit				-	-										16.10	23.45	13.15		
Earth—Elliptical Orbit .		-			Promise								_		22.70	30.10	19.80		
Orbit-Elliptical Orbit .														-	13.25	20.60	10.30		

¹ No return trip included.

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¹ No return trip included

Earth—Elliptical Orbit. Orbit—Elliptical Orbit.

16.10 23.45 13.15 22.70 30.10 19.80 13.25 20.60 10.30

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